

Multi-scale Modeling using the Adaptive Resolution Scheme

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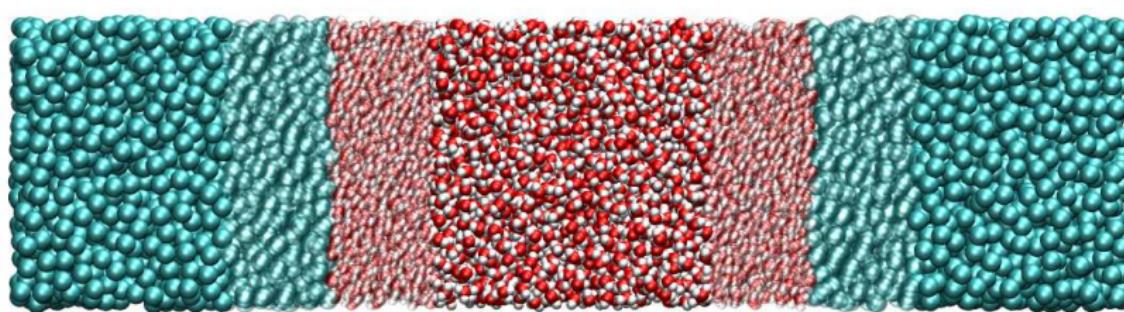
Coworkers

A lot of people have contributed to AdResS:

- Cecilia Clementi (Rice University, Texas)
- Luigi Delle Site (MPI-P → FU Berlin)
- Sebastian Fritsch (MPI-P)
- Kurt Kremer (MPI-P)
- Brad Lambeth (Rice University, Texas)
- Debashish Mukherji (MPI-P)
- Patrick Kiley (University of Cambridge)
- Simón Poblete (MPI-P → FZ Jülich)
- Adolfo Poma (MPI-P → Sapienza University of Rome)
- Matej Praprotnik (MPI-P → National Institute of Chemistry, Slovenia)
- Stac Bevc (National Institute of Chemistry, Slovenia)

Introduction

The Adaptive Resolution Scheme (AdResS)



The big idea¹

- Coupling of two resolutions (AA and CG) by a force interpolation:

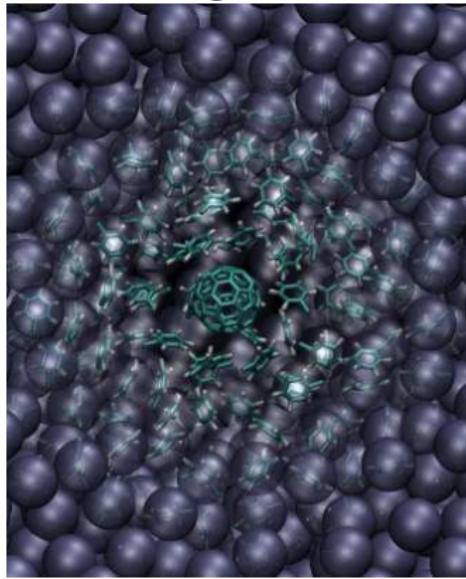
$$\vec{F}_{\alpha\beta} = w_\alpha w_\beta \vec{F}_{\alpha\beta}^{\text{AA}} + [1 - w_\alpha w_\beta] \vec{F}_{\alpha\beta}^{\text{CG}} ,$$

where α and β are molecules

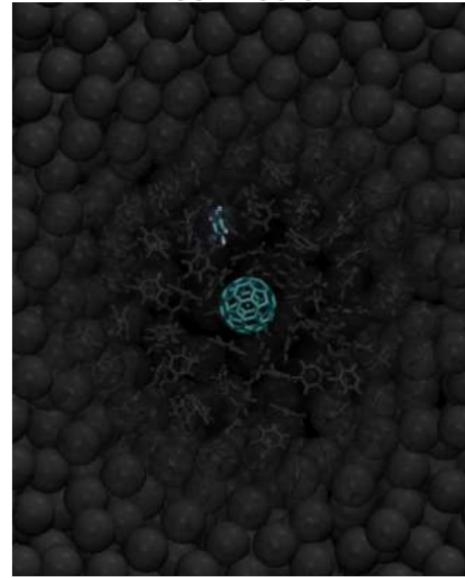
¹M. Praprotnik, L. Delle Site, and K. Kremer, JCP **123**, 224106 (2005).

Introduction Illustration

Magnifier:

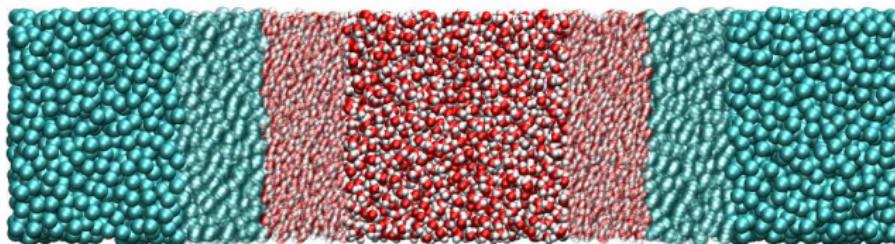


Free motion:



Introduction

Motivation



Inherited from coarse-graining

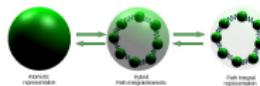
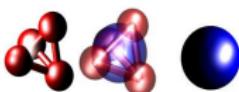
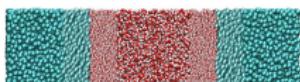
- Accessibility of bigger time- and length-scales
 - Avoids sampling of unimportant degrees of freedom
 - Better sampling through smoother energy landscape
 - Computational speed-up

As an analysis tool

- Study influence of coarse-graining
 - Replace surrounding environment

Introduction

What kind of systems can be coupled?



- Atomistic and coarse-grained representation of the same system
 - Atomistic and very simplified (LJ) system
 - Quantum and atomistic systems

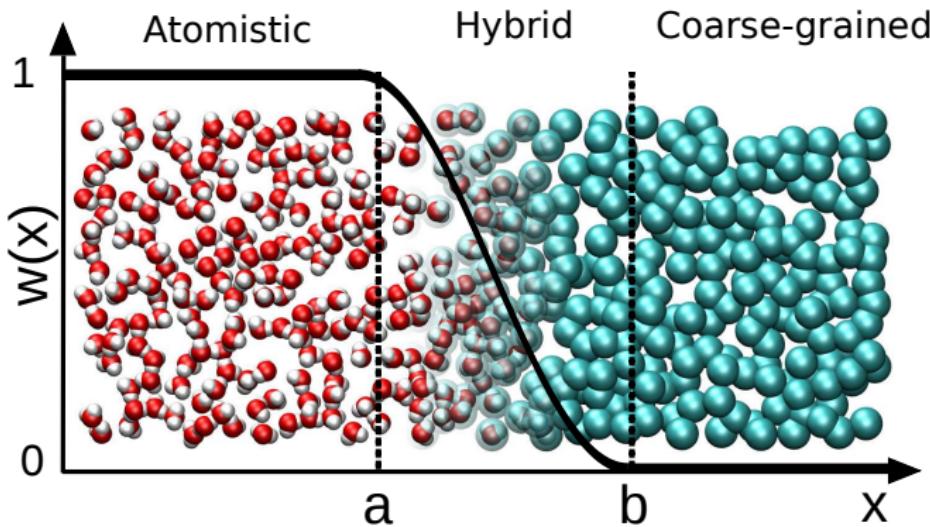
Coupling conditions²

$$\rho^{\text{at}} = \rho^{\text{cg}}, \quad T^{\text{at}} = T^{\text{cg}}, \quad (p^{\text{at}} = p^{\text{cg}})$$

²M. Praprotnik, L. Delle Site and K. Kremer; Annu. Rev. Phys. Chem. **59**, 545 (2008).

Methodic Details

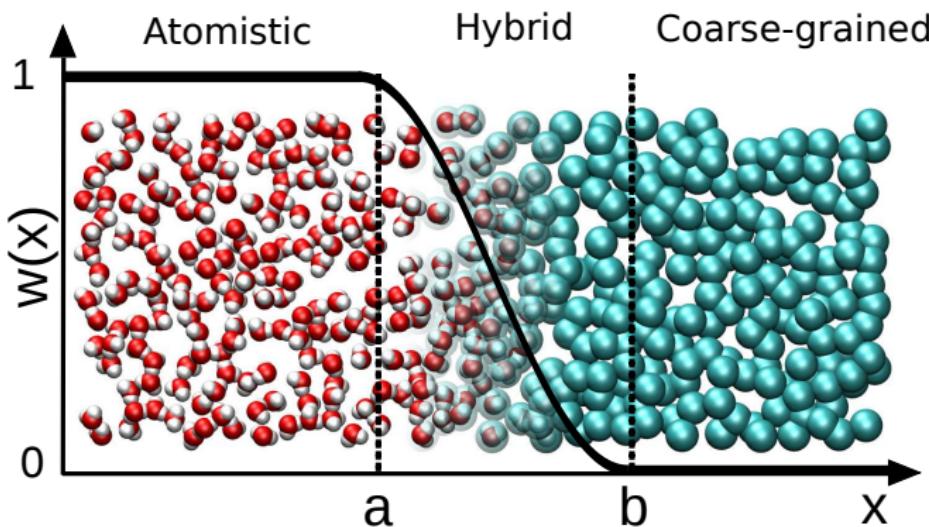
The Weighting Function I



$$w(x) = \begin{cases} 1 & : \text{atomistic/explicit region} \\ 0 < w < 1 & : \text{hybrid region} \\ 0 & : \text{coarse-grained region} \end{cases}$$

Methodic Details

The Weighting Function II

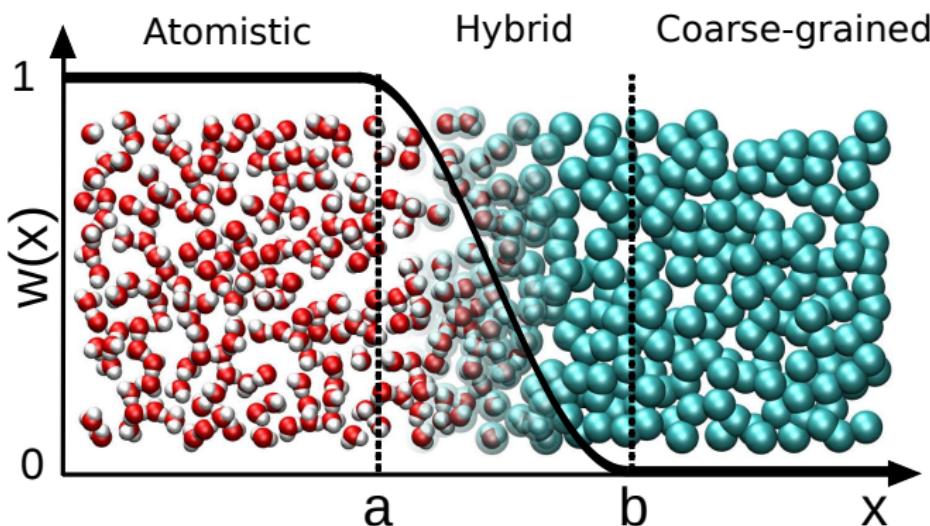


Boundary conditions

- Monotonic \rightarrow No artificial barriers
- $w'(d_{\text{ex}}) = 0 = w'(d_{\text{ex}} + d_{\text{hy}}) \rightarrow$ smooth free energy transition

Methodic Details

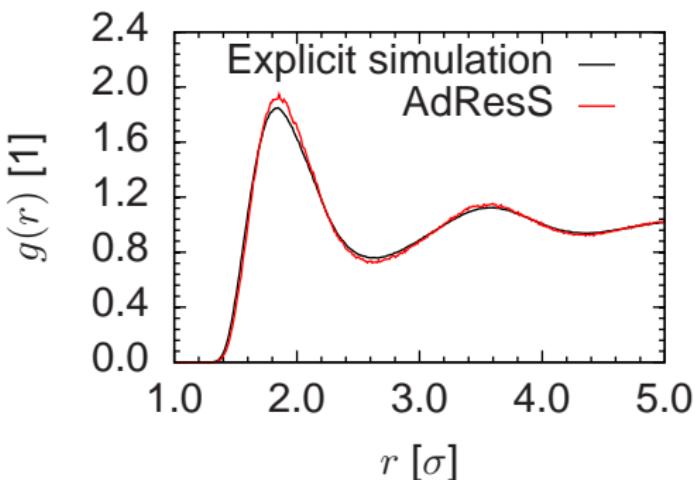
The Weighting Function III



$$w(x) = \begin{cases} 0 & : x > d_{\text{ex}} + d_{\text{hy}} \\ \cos^2 \left(\frac{\pi}{2d_{\text{hy}}} (x - d_{\text{ex}}) \right) & : d_{\text{ex}} + d_{\text{hy}} > x > d_{\text{ex}} \\ 1 & : d_{\text{ex}} > x \end{cases}$$

AdResS at a First Glance

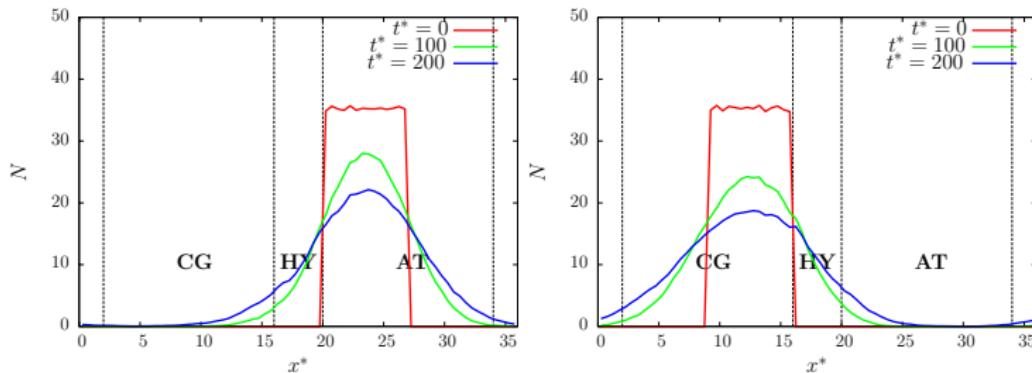
Radial Distribution Function (Tetrahedral Liquid)



- AdResS as a “magnifier” for a certain region
- RDF matches and accuracy depends more on the accuracy of the coarse-graining than on AdResS

AdResS at a First Glance

Particle Fluctuations

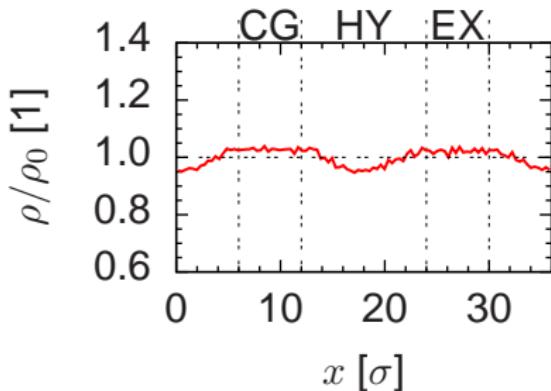


- Particles can freely diffuse through the hybrid region
- Profile is slightly asymmetric
- Coarse-grained dynamics is slightly faster

AdResS at a First Glance

Density Artifacts in the Hybrid Region

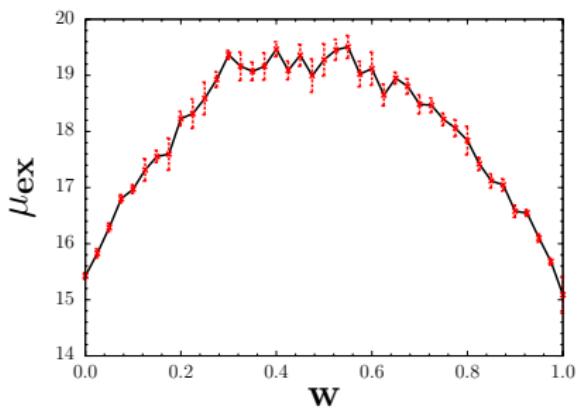
There is no free lunch!



- System is inhomogeneous in the hybrid region
- Linear interpolation is too simple
- Can be corrected, but is not necessary in all cases

Theory of the Thermodynamic Force

Idea



- Chemical potential in the hybrid region is different
- Add an external field³

$$\vec{f}_{\text{thermo}} = -\nabla(\mu_{AA} - \mu(w))$$

to the linear interpolation to repair the density inhomogeneity

- Correction is not always necessary
- Allows to couple systems of different pressures

³S. Poblete, M. Praprotnik, K. Kremer and L. Delle Site, JCP **132**, 114101
(2010)

Theory of the Thermodynamic Force

Pressure View

- Chemical potential is tedious to measure
- Assumption of constant weight
- To obtain a flat density profile one needs to compensate for the different pressures in the hybrid region⁴

$$\vec{f}_{\text{thermo}} = -\frac{1}{\rho_0} \nabla(p_{AA} - p(w))$$

- Local pressure depends on the local density

$$\vec{f}_{\text{thermo}} = -\frac{1}{\rho_0^2 \kappa} \nabla(\rho(w))$$

- Easier, because $\rho(w)$ can be measured on the fly
- Iterative correction is possible

⁴S. Poblete, S. Fritsch, C. Junghans, G. Cicotti, K. Kremer and L. Delle Site, PRL **108**, 170602 (2012)

Numerical Details

Overview I

The problem is the hybrid region!

Question

- How to reintroduce the atomistic details?
- How to initialize the atomistic velocities?
- How to introduced the correct bond distribution?
- How to map the molecular force interpolation to the atoms in the atomistic region?

Numerical Details

Overview II

Consumed Computer Times in MD

- 70% Force calculation
- 15% Communication
- 8% Neighbor search
- 5% Position update
- 2% IO

Position update is cheap → do NOT remove atomistic particles, but do not treat them force-wise and integrate them everywhere.

Double representation trick⁵ ( everywhere)

⁵C. Junghans and S. Poblete, Comp. Phys. Comm. **181**, 1449 (2010).

Numerical Details

Virtual Sites

How to implement the double representation?

Virtual sites: Dummy particles, which are moved according to a geometrical rule rather than by Newton's equation of motion.



Force distribution

Force on an atom in a molecule with a virtual site:

$$\vec{F}_i^{\text{atom}} = -\frac{\partial(V + V^{\text{cg}})}{\partial \vec{r}_i} = \vec{F}_i^{\text{ex}} + \vec{F}^{\text{cg}} \frac{\partial \vec{r}^{\text{cg}}}{\partial \vec{r}_i} \stackrel{\text{COM}}{=} \vec{F}_i^{\text{ex}} + \frac{m_i}{\sum_{i \in \alpha} m_i} \vec{F}^{\text{cg}}$$

Can be easily added to the integrator!

Numerical Details

Integrator

Velocity Verlet Integrator (with AdResS and virtual sites)

1. $v(t + \Delta t/2) = v(t) + \Delta t/2 f(t)/m$
2. $p(t + \Delta t) = p(t) + \Delta t v(t + \Delta t/2)$
- 2b. Update the positions, velocities and weighting functions $w(\vec{R})$ of the virtual sites
3. Calculate $f(t + \Delta t)$ from $p(t + \Delta t)$, $v(t + \Delta t/2)$ (and thermostat)
- 3b. Distribute the force of the virtual sites to the real particles
4. $v(t + \Delta t) = v(t + \Delta t/2) + \Delta t/2 f(t + \Delta t)$

Virtual sites increase the amount of communication ($\mathcal{O}(N_{\text{VS}}^{2/3})$).

Numerical Details

Thermostat

A thermostat is needed for a smooth transition.

Langevin Thermostat

Thermostat force acting on the atoms

$$\vec{f}_i^{\text{thermo}} = \vec{f}_i^D + \vec{f}_i^R = \xi_i \vec{v}_i + \sigma_i \vec{\eta}_i .$$

One can show that

$$\sum_{i \in \alpha} \vec{f}_i^{\text{thermo}} = \vec{F}_{\alpha}^D + \vec{F}_{\alpha}^R$$

is also a Langevin thermostat with a common fluctuation-dissipation theorem ($\sigma^2 = k_B T \xi$) if the friction constants ξ_i have been scaled with the masses, $\xi_i = m_i \xi$.

DPD thermostat can also be used, but more theory is involved.

Numerical Details

Reinitialization

Possibilities for Velocity and Bond Reinitialization

- Initialization from Maxwell/bond distribution
- Freeze and unfreeze them
- Keep on integrating them

Reminder: We have hybrid molecules everywhere
Bonded interactions are cheap
→ keep on integrating them is the easiest solution

Numerical Details

Cut-Offs

AdResS needs a molecule-based cut-off due to the force interpolation.

Molecular Cut-Offs

- Avoid that molecules interact half atomistic - half coarse-grained
- Avoid dipoles at the cut-off

Numerical Details

Speed-up

Simulation of a big system with small atomistic region: ⁶

Ideal strong speed-up (const. systemsize)

$$\begin{aligned} s^{-1} &= \left(\frac{t_{\text{new}}}{t_{\text{old}}} \right) \\ V_{\text{AA}} \ll V_{\text{CG}} &\quad \frac{t_{\text{not-forces}}}{t_{\text{total}}} + \frac{t_{\text{forces}}}{t_{\text{total}}} \frac{1}{(\text{atoms per molecule})^2} \\ \text{Toluene} &\equiv 0.26 + 0.74 \cdot \frac{1}{2 \times 12^2} \approx 3.6^{-1} \end{aligned}$$

⁶S. Fritsch, C. Junghans and K. Kremer, JCTC **8**, 398 (2012).

Basic Examples

Tetrahedral Liquid

Toy model for a liquid (methane-like). 7



Explicit (all-atom) Interactions

Non-bonded (even inside the molecules), LJ:

$$U_{LJ}^{ex}(r_{i\alpha j\beta}) = \begin{cases} 4\epsilon \left[(\sigma/r_{i\alpha j\beta})^{12} - (\sigma/r_{i\alpha j\beta})^6 + \frac{1}{4} \right], & r_{i\alpha j\beta} \leq 2^{1/6}\sigma \\ 0, & r_{i\alpha j\beta} > 2^{1/6}\sigma \end{cases}$$

Bonded, FENE:

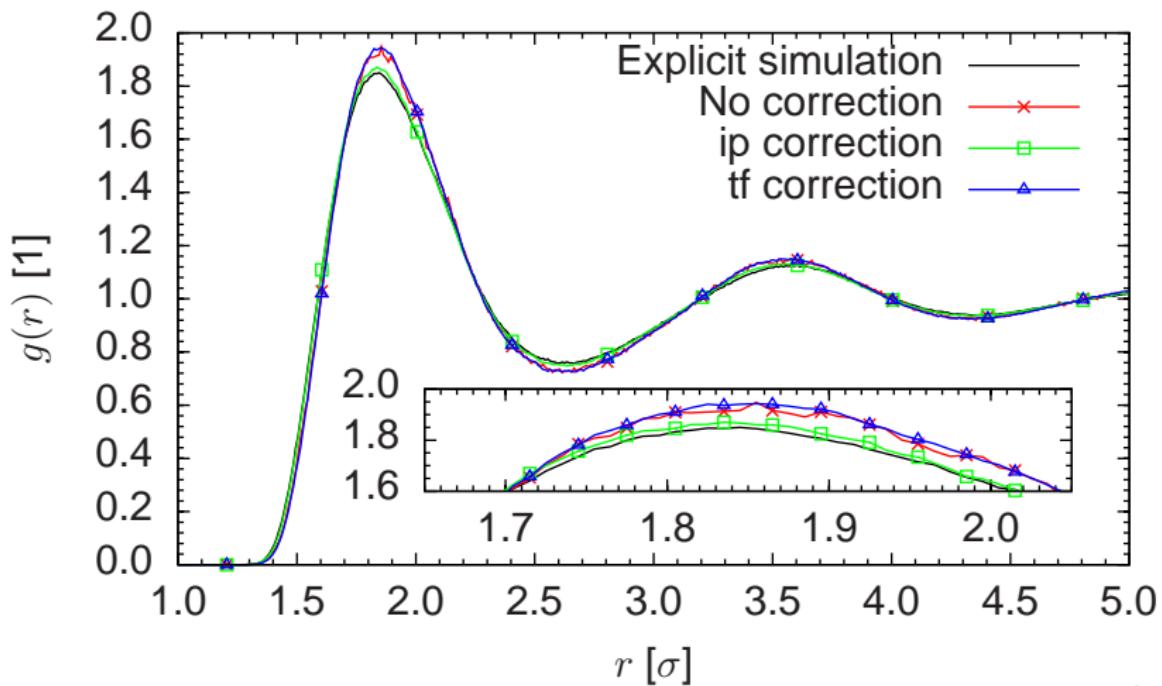
$$U_{FENE}^{ex}(r_{i\alpha j\alpha}) = \begin{cases} -\frac{1}{2}kR_0^2 \ln \left[1 - (r_{i\alpha j\beta}/R_0)^2 \right], & r_{i\alpha j\beta} \leq R_0 \\ \infty & r_{i\alpha j\beta} > R_0 \end{cases}$$

Coarse-grained Interactions

Isotropic one-site tabulated potential $U(r_{ij})$ by IBI

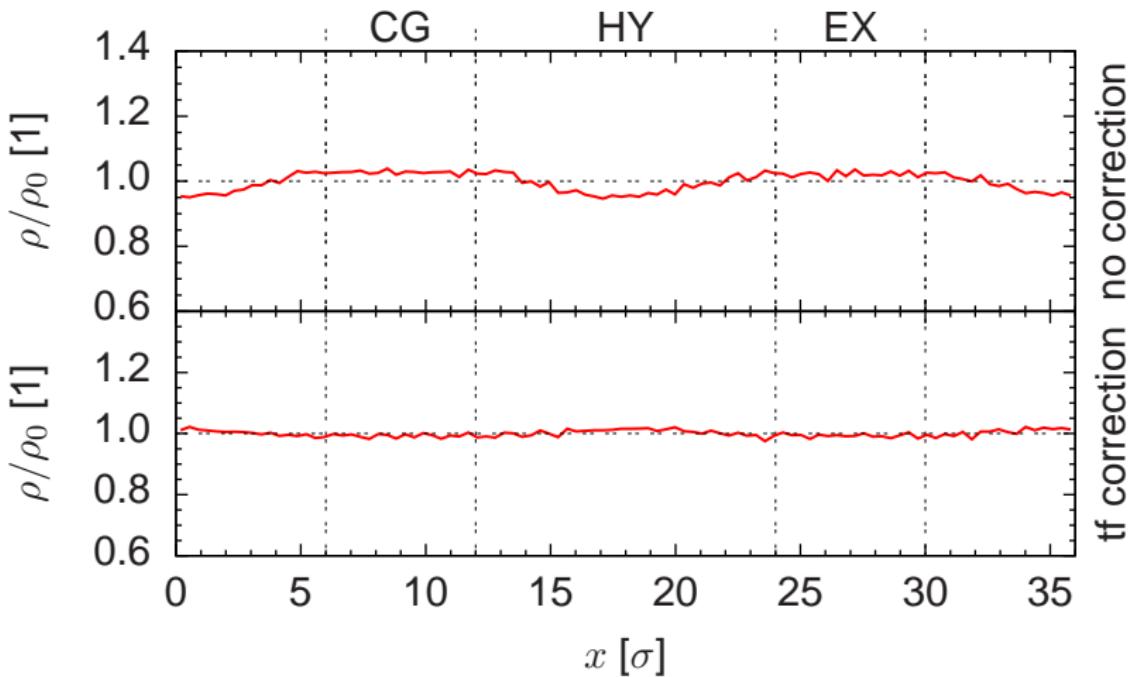
Basic Examples

Tetrahedral Liquid (Structure)



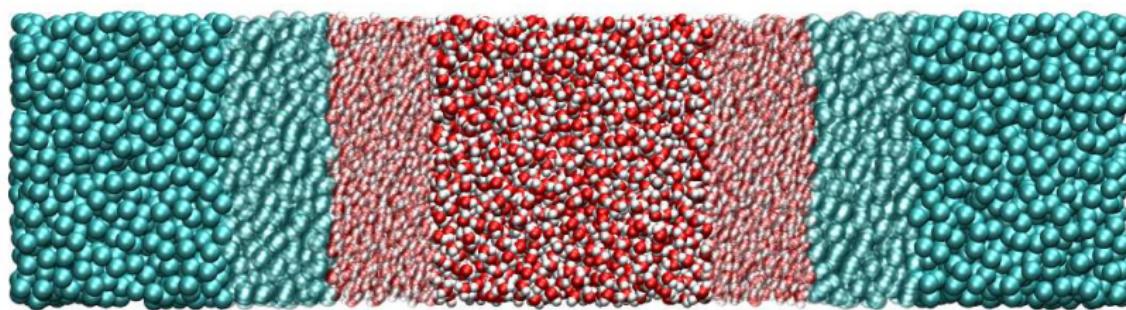
Basic Examples

Tetrahedral Liquid (Density profile)



Basic Examples

Water



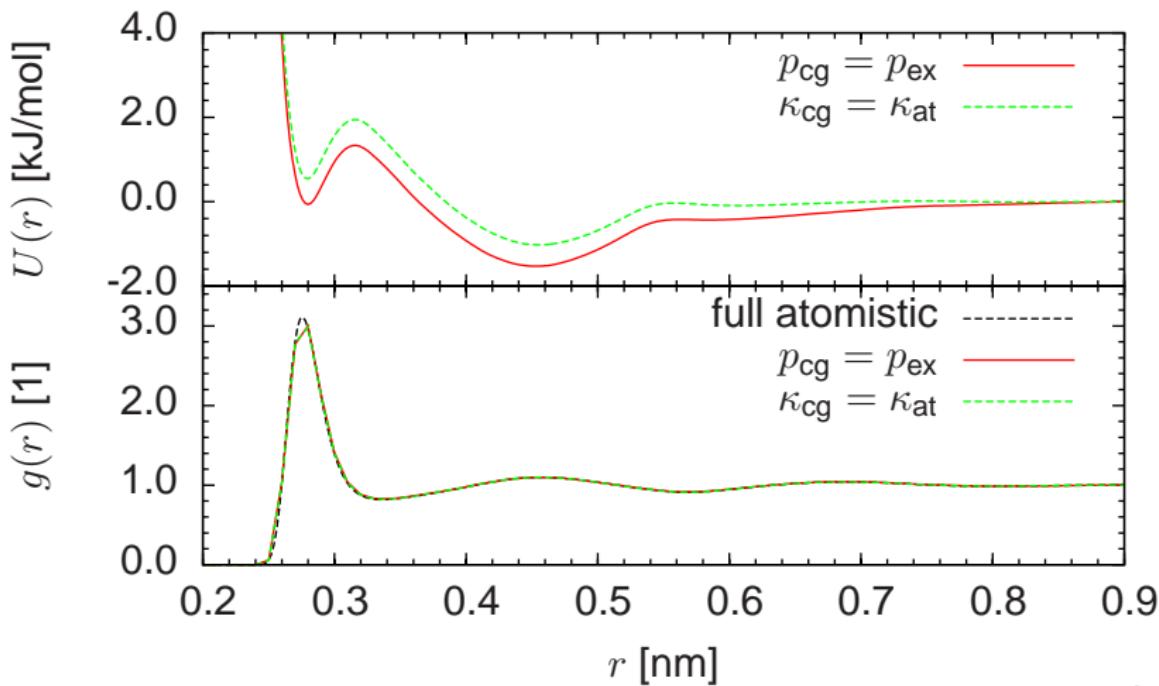
- A realistic, but simple system
- Electrostatics by Reaction Field
- One bead coarse-grained model

Can we couple a not pressure-corrected model, too? ⁸

⁸S. Poblete, S. Fritsch, C. Junghans, G. Cicotti, K. Kremer and L. Delle Site, PRL **108**, 170602 (2012)

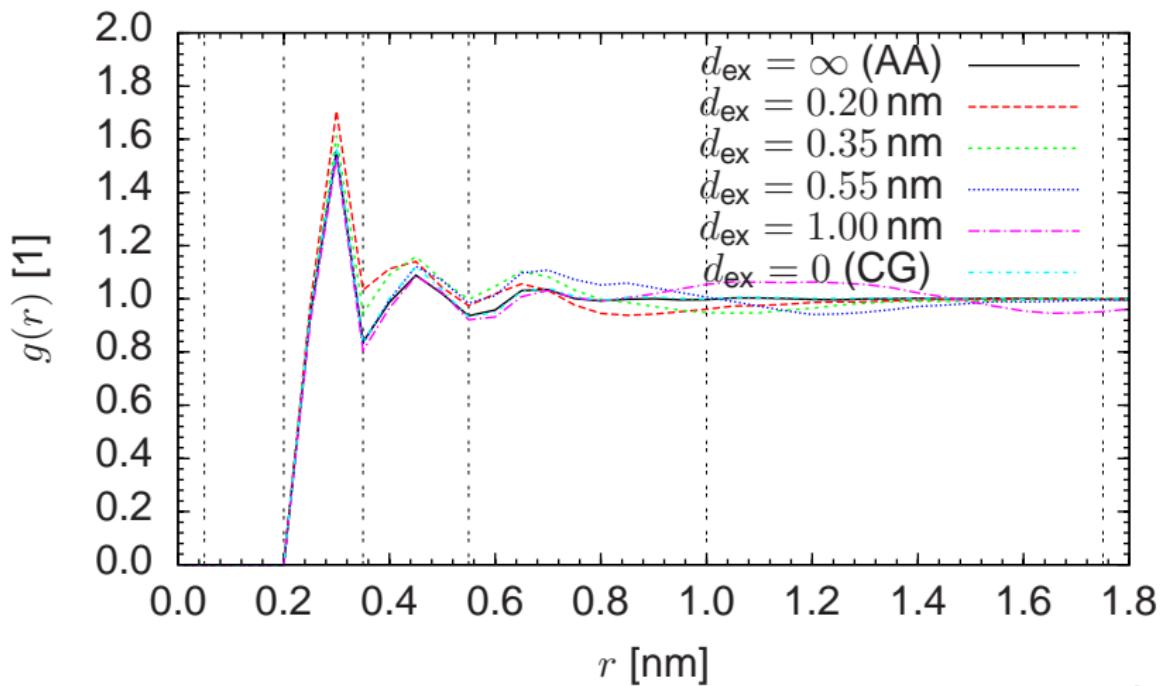
Basic Examples

Water (Coarse-grained Model)



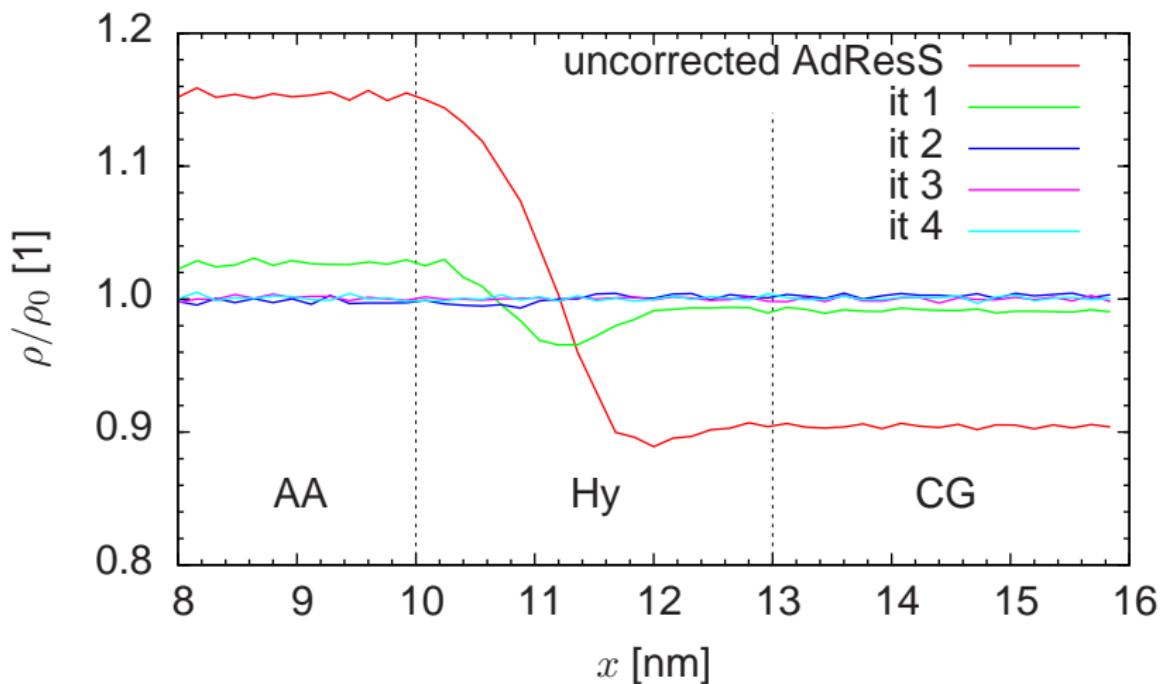
Basic Examples

Water (Perturbation without Thermodynamic Force)



Basic Examples

Water (Density profile)

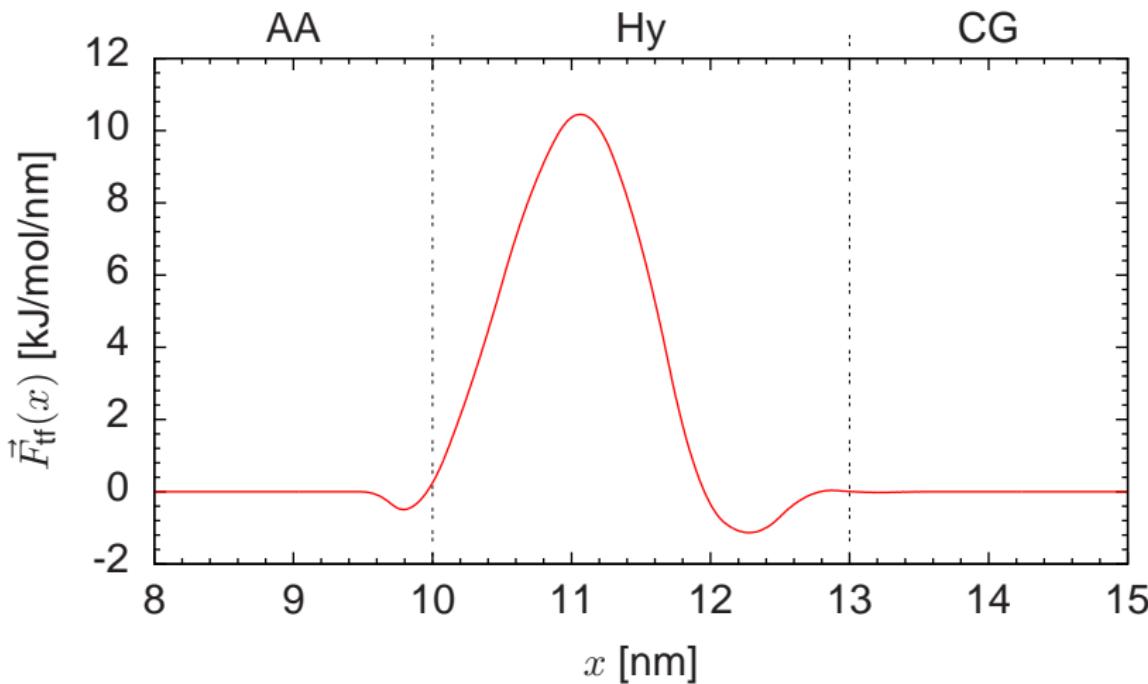


One can even bring non-pressure corrected model in equilibrium



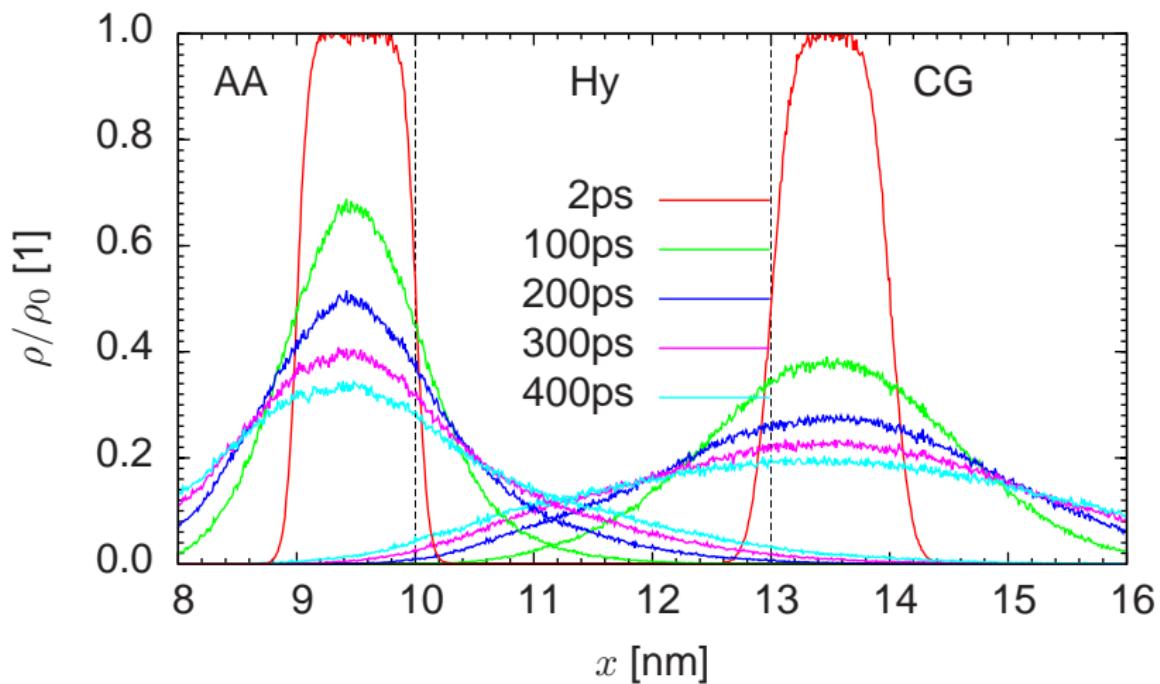
Basic Examples

Water (Thermodynamic Force)



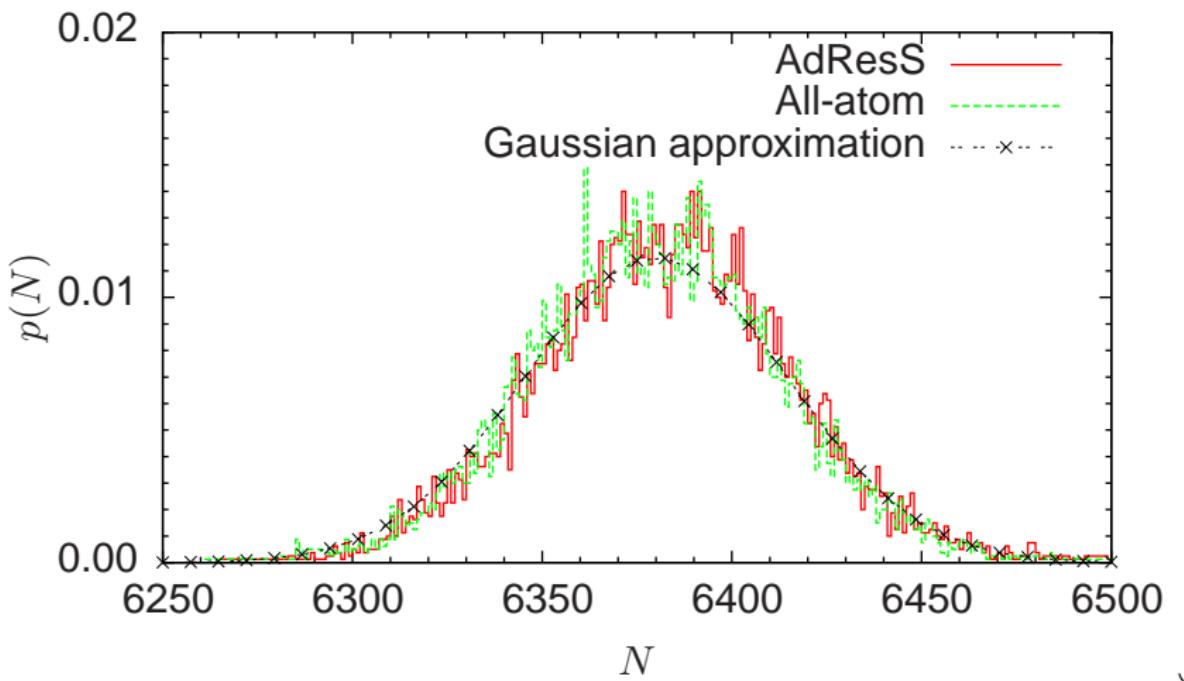
Basic Examples

Water (Diffusion Profile)



Basic Examples

Water (Particle Fluctuations in AA zone)

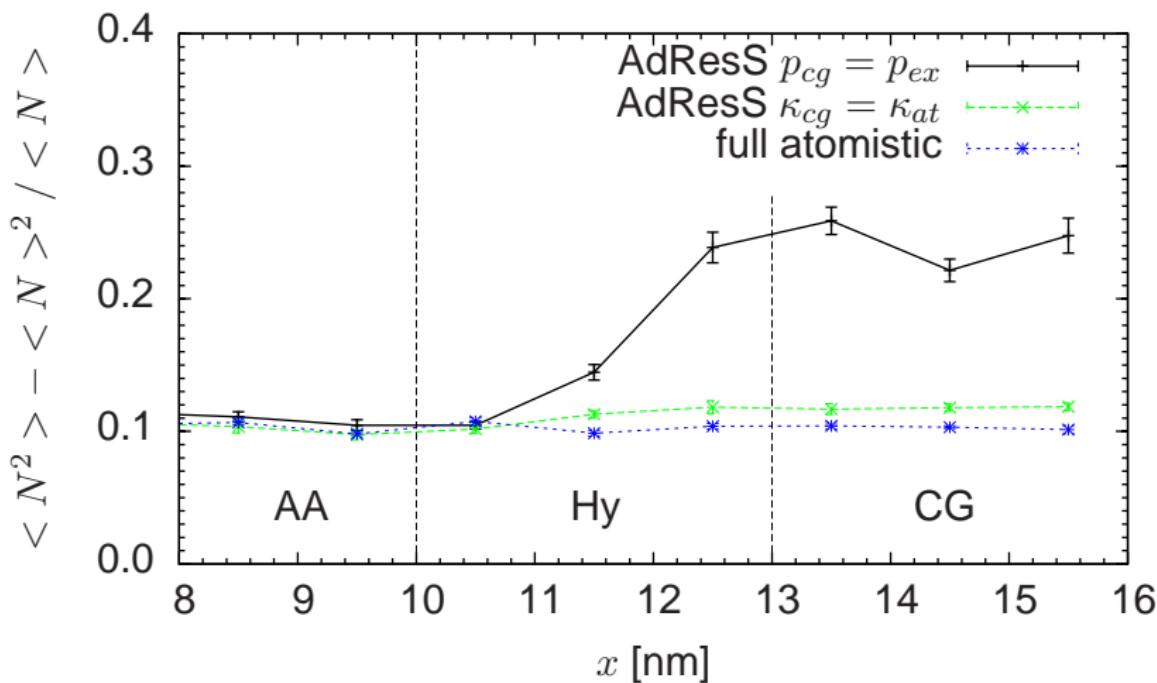


Grand-canonical behaviour



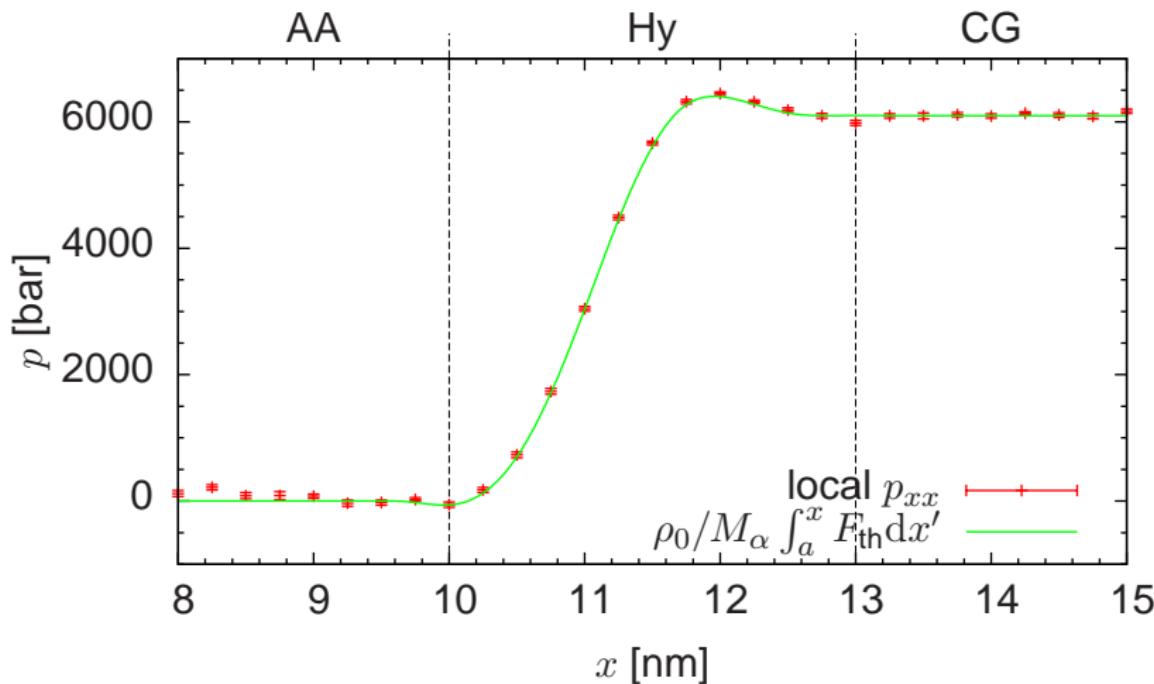
Basic Examples

Water (Particle Fluctuations)



Basic Examples

Water (Thermodynamic Force)



Thermodynamic force compensates the pressure difference.

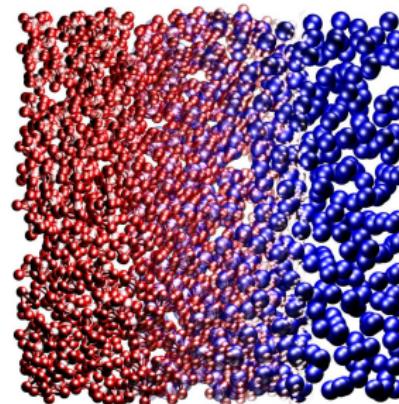
Outlook

AdResS:

- Useful analysis tool
- Speed-up simulations
- Molecules exchange between the two resolutions

Things to work on:

- Electrostatics
- Adaptive implementation
→ ESPResSo ++ ✓



The End

Thank you for your attention !

